LIDAR, AERONET AND AEROSOL-CLOUDS INTERACTIONS

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INTRODUCTION WHAT IS REMOTE SENSING?







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Process of obtaining information about properties of a target without contact.

This information is propagated by means of electromagnetic radiation interacting with the target.



Lidar (light detection and ranging) is an active remote sensing technology that measures distance by illuminating a target with a laser and analyzing the reflected light

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There are a wide variety of fields of applications for lidar:

- * Geology
- * Archaeology
- * Biology (vegetation)
- * Architecture (cultural heritage)
- * Astronomy
- * Atmosphere



Geology: In geophysics and tectonics, a combination of aircraft-based lidar and GPS has evolved into an important tool for detecting faults and for measuring uplift. The output of the two technologies can produce extremely accurate elevation models for terrain and enables surveys to be taken of the coastline. It is also used extensively to monitor glaciers and perform coastal change analysis





Archaeology: aiding in the planning of field campaigns, mapping features beneath forest canopy, and providing an overview of broad, continuous features that may be indistinguishable on the ground, creating high-resolution digital elevation models (DEMs) of archaeological sites that can reveal micro-topography that are otherwise hidden by vegetation



Image taken from: Chase et al., 2011, Journal of Archeological Science, 38, 387-398 Maya Site of Caracol (Belize)



Biology (vegetation): lidar has also found many applications in forestry. Canopy height and biomass measurements can be all be studied using airborne lidar systems. Tree height strongly correlated to other properties (as woody biomass, age, etc)

LASER-SCANNING





Architecture (cultural heritage): fluorescence lidar technique has been applied to the investigation of the cultural heritage only quite recently: remote diagnostics of monuments, providing helpful information for the assessment of the conservation status of monuments in the outdoor and for the characterization of the materials employed in their construction. Main applications include the detection and characterization of different stones, mortars and other construction materials, of protective treatments, of biodeteriogens and the study of the effects of biocide treatments.



Figure 3.6.3: Basic principles of lidar hyper-spectral imaging on monuments [17].

© Raimondi et al., in: "Handbook on the Use of Lasers in Conservation and Conservation Science", 2008.

- Aerosol particles: liquid or solid particles suspended in a gaseous medium (several types)
- Atmospheric aerosols play an important role in the Earth's climate system:
 - 1. they interact with solar and thermal radiation, modulating the Earth radiation budget
 - 2. they modify clouds microphysical properties by acting as cloud condensation nuclei and ice forming nuclei



Aerosol particles: cooling/warming effect

depending on vertical distribution and types (among other factors)

- Classification:
 - ✓ Mineral particles
 - ✓ Marine particles
 - ✓ Carbon particles
 - ✓ Sulfates, nitrates and organic compounds



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Radiative forcing by components





IPCC (2013)







LIDAR TECHNIQUE



The key of this technique is the use of **pulsed lasers**, which allow for obtaining range-resolved information from the delay between emitted and received pulses

RANGING (BASED ON TIMING)



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RANGING (BASED ON TIMING)

$$\Delta R = R_1 - R_2 = \frac{c\tau}{2}$$

"effective (spatial) pulse length"

(**potential** spatial lidar resolution)

Spatial resolution example:

Nd:YAG Laser (1064 nm) Q-switch pulsed: τ ~ 10 ns

 $\Delta R \approx 1.5 \text{ m}$

In practice there are additional limitations like the sampling frequency of the acquisition system ...



LIDAR SETUP



- Transmitter system \succ
- Receiver system



LIDAR SETUP: LASER SOURCE

- ➤ 1960's: nitrogen laser (337 nm) and ruby laser (694nm→347nm): relatively low average power, limited to 2 km
- 1980's: pulsed lasers with high average power as UV excimer laser (XeCl at 308nm and XeF at 351 nm) and Nd:YAG laser
- ➤ currently Nd:YAG (Neodymiun-doped Yttrium Aluminum Garnet): 1064 nm → 532 nm → 355 nm repetition frequency : 1-50 Hz energy/pulse: 100 mJ to 1.5 J (at 1064 nm)





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LIDAR SETUP: TRANSMITTER OPTICS

> Mirrors to transmit the laser beam (high reflectivity / high transmissivity)

Beam expander made of lenses (anti-reflection coated) allows for: a beam expansion by a factor ~ x4 - x10 reducing background light increasing SNR reducing divergence (~1mrad)



LIDAR SETUP: RECEIVER OPTICS

φ_{primary}: 30 cm – 1 m and φ_{secondary}: ~ 1cm
compromise between a small FOV necessary for high background supression and a large FOV for a sufficient signal intensity from short distances (FOV a factor of ~2-10 larger than the laser divergence)
receiver optics behind telescope must be optimized for high transmission of the Raman signals

 \triangleright dichroic beam splitters: reflect light of a certain λ an transmit others

▶ interference filters: typically FWHM <0.5 nm, suppression factor 10⁸-10¹⁰



Telescopes





Newtonian



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LIDAR SETUP: DETECTORS AND DATA ACQUISITION

- photomultipliers in An and PC mode are typically used (PMTs)
- ▶ high quantum efficiency (in UV ~ 25%) and low noise
- detector outputs can be preamplified before registration
- \succ time resolution (or window length) is ~ 100 ns, (7.5 15 m)
- \blacktriangleright averaging time for the raw signals is 10- 60 s
- ➤ signals are usually further averaged in time and space during data evaluation



ELASTIC LIDAR FOR AEROSOLS

Generally consists of a non-tunable high-power pulsed laser and intensity of the received signal is measured

A narrow beam is transmitted into the atmosphere and backscattered by the atmosphere to a receiver telescope and detector

The nature of the backscattering is determined by the properties of the volume of the atmosphere that contains the Rayleigh (molecules) and Mie (particles) scatterers

The combination of the short laser pulse (~10 ns) and the small beam divergence (~ 10^{-3} to 10^{-4} radians) results in volumes of a few cubic meters at ranges of tens of km

The primary properties measured are the intensity and polarization of the signal, and these are used to retrieve particle properties



LIDAR EQUATION



P_o average power of single laser pulse, τ temporal pulse length E_o=P_oτ pulse energy The propagation through distance R implies an attenuation $e^{-\int_{0}^{R} \alpha(x) dx}$ and $(P_o / A_L) e^{-\int_{0}^{\alpha(x) dx}}$ represent the power per unit area that illuminates a volume of atmosphere at distance R

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The illuminated volume $c(\tau/2)A_L$ produces a backscattered signal proportional to the backscatter coefficient This signal travels back to the detector suffering an additional attenuation and thus the signal collected by the detector system of surface A is:



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LIDAR EQUATION

$$P(R,\lambda) = P_o(\lambda) \frac{C}{R^2} \beta(R,\lambda) e^{-2\int_0^R \alpha(x,\lambda) dx} = P_o(\lambda) \frac{C}{R^2} \beta(R,\lambda) T(R,\lambda)$$

 $P(R,\lambda)$ = power received by the system after backscattering from range R

 $P_o(\lambda)$ = power transmitted

C = lidar calibration constant

 $\beta(R, \lambda)$ = backscatter coeff. (length⁻¹sr⁻¹)

 $\alpha(R, \lambda)$ = volume extinction coeff. (length⁻¹)

The term $\beta(R, \lambda)$ is the backscatter coefficient at distance R, it stands for the ability of the atmosphere to scatter light back into the direction from which it comes

T(R, λ) is the transmission term and describes how much light gets lost on the way from the lidar to distance R and back

Most of the information about atmospheric properties derived from backscatter lidar measurements is based on extinction, $\alpha(R, \lambda)$, and backscatter coefficients, $\beta(R, \lambda)$

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... MORE ON LIDAR EQUATION

Overlap function O(R): geometrical overlap between laser beam and telescope field of view

O(R) accounts for the partial overlap in the near height-range and tends to estabilize (ideally to 1) in the far height-range



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Elastic lidar equation including overlap:

$$P(R,\lambda) = P_o(\lambda) \frac{C}{R^2} O(R) \beta(R,\lambda) e^{-2\int_0^R \alpha(x,\lambda) dx}$$



... MORE ON LIDAR EQUATION



$$P(R,\lambda) = P_o(\lambda) \frac{C}{R^2} O(R) \beta(R,\lambda) e^{-2 \int_0^R \alpha(x,\lambda) dx}$$





 $\beta(R, \lambda)$ = backscatter coeff. (length⁻¹sr⁻¹) $\alpha(R, \lambda)$ = volume extinction coeff. (length⁻¹)

Depend on the particle properties as size, shape, composition, concentration, etc.



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FIRST STEPS USING LIDAR EQUATION

Signal-to-noise ratio (SNR) is a critical issue when atmospheric properties are computed. To improve SNR, two methods are possible:



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FIRST STEPS USING LIDAR EQUATION

...performing a basic calculation (range corrected signal) on the lidar equation, a first idea of the behavior of the atmosphere can be obtained...



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RECOMMENDED BIBLIOGRAPH

KOVALEV, V.A, EICHINGER, W.E. Elastic Lidar. Wiley Interscience, New Jersey, 2004. Chapters 3 and 4.

WEITKAMP, C., Lidar. Range-resolved optical remote sensing of the Atmosphere, Springer, New York, 2005. Chapter 1.







GEOMETRICAL PROPERTIES

Geometrical properties:

- base height -
- top height
- geometrical thickness
- center of mass



Quicklooks can provide а rough estimate of geometrical properties but quantitative information is needed



GEOMETRICAL PROPERTIES

To quantify geometrical properties two groups of algorithms: threshold methods and derivative methods

Threshold methods: processed profiles are needed. Typically, this threshold is established in terms of the backscattering ratio:

$$\begin{aligned} R_{back}(R) &= \frac{\beta_{part}(R) + \beta_{mol}(R)}{\beta_{mol}(R)} = \\ &= 1 + \frac{\beta_{part}(R)}{\beta_{mol}(R)} \end{aligned}$$



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Derivative methods:

- advantage of the derivative methods: applied directly to uncalibrated raw lidar data, i.e., no additional measurements, models or assumptions are required
- first derivate method, second derivate method, logarithm derivate method
- logarithmic derivate method because it mostly quantifies the relative change of the signal, instead of an absolute change as other derivative methods

- the backscattered lidar signal of a clear atmosphere decreases monotonically with altitude (derivative is always negative)
- aerosol layer implies an abrupt increase in the signal values, such that around the layer base and layer top the sign of this derivative changes
- the change from negative to positive is identified as layer boundaries

















LIDAR OPTICAL ALGORITHMS

First we will describe the technique used to compute the particle backscatter coefficient from return signals measured with the widely used elastic standard backscatter lidar

The main drawback of this method is that trustworthy profiles of the climate relevant volume extinction coefficient of the particles cannot be obtained. The extinction profile must be estimated from the determined backscatter coefficient profile

RAMAN METHOD

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By applying the so-called Raman lidar technique, the profile of the particle extinction and backscatter coefficients can be independently determined

An aerosol Raman lidar measures two signal profiles, which permit the separation of particle and molecular backscatter contributions

Starting point: lidar equation

$$P(R,\lambda) = P_o(\lambda) \frac{C}{R^2} O(R) \beta(R,\lambda) e^{-2\int_0^R \alpha(x,\lambda) dx} =$$
$$= P_o(\lambda) \frac{C}{R^2} O(R) \beta(R,\lambda) T(R,\lambda)^2$$

 $\beta(R,\lambda)$ (km⁻¹sr⁻¹) and $\alpha(R, \lambda)$ (km⁻¹) are the backscatter and extinction coefficients in the atmosphere

Remembering the definition of lidar range corrected signal:

$$R.C.S.(R,\lambda) \equiv P(R,\lambda) \cdot R^2 = P_o(\lambda) \cdot C \cdot O(R) \cdot \beta(R,\lambda) \cdot e^{-2\int_0^{\Lambda} \alpha(x,\lambda) dx}$$



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 $R.C.S.(R,\lambda) \equiv P(R,\lambda) \cdot R^2 = P_o(\lambda) \cdot C \cdot O(R) \cdot \beta(R,\lambda) \cdot e^{-2\int_0^R \alpha(x,\lambda) dx}$

The conversion of lidar signal into range corrected lidar signal causes a change in appearance:



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KLETT-FERNALD METHOD

$$R.C.S.(R,\lambda) \equiv P(R,\lambda) \cdot R^2 = P_o(\lambda) \cdot C \cdot O(R) \cdot \beta(R,\lambda) \cdot e^{-2\int_0^{\alpha(x,\lambda)dx} dx}$$

... some considerations:

→ the overlap is assumed to be complete , O(R) =1 → $R \ge R_{min}$

► we dropp the wavelength dependence $R.C.S.(R) \equiv P(R) \cdot R^2 = P_o \cdot C \cdot \beta(R) \cdot e^{-2\int_0^R \alpha(x)dx}$

> $\beta(R)$ (km⁻¹sr⁻¹) and $\alpha(R)$ (km⁻¹) are caused by particles and molecules:

$$\beta(R) = \beta_{mol}(R) + \beta_{part}(R)$$

$$\alpha(R) = \sigma_{mol}^{scat}(R) + \sigma_{mot}^{abs}(R) + \sigma_{part}^{scat}(R) + \sigma_{part}^{abs}(R)$$

 $\approx 0 \text{ km}^{-1}$ at lidar wavelengths

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Taking into account molecular absorption effects are ignored:

$$\alpha(R) = \sigma_{mol}^{scat}(R) + \sigma_{part}^{scat}(R) + \sigma_{part}^{abs}(R) = \alpha_{mol}(R) + \alpha_{part}(R)$$

Therefore, the range corrected signal is:

$$\underline{R.C.S.(R)} = P_o \cdot C \cdot \left[\beta_{mol}(R) + \beta_{part}(R)\right] \cdot e^{-2\int_0^R \left[\alpha_{mol}(x) + \alpha_{part}(x)\right] dx}$$

unknowns

measurement

The molecular properties, $\beta_{mol}(R)$ and $\alpha_{mol}(R)$, can be determined from the Rayleigh theory using the best available meteorological data of temperature and pressure or approximated from appropriate standard atmospheres so that only the aerosol scattering and absorption properties, $\beta_{part}(R)$ and $\alpha_{part}(R)$, remain to be determined

Shortcoming: 1 measurement versus 2 unknowns

$$\underbrace{R.C.S.(R)}_{\text{measurement}} = P_o \cdot C \cdot \left[\beta_{mol}(R) + \beta_{part}(R)\right] \cdot e^{-2\int_0^R \left[\alpha_{mol}(x) + \alpha_{part}(x)\right] dx}$$
unknowns

How to solve?: definition of the particle and molecular lidar ratio (extinction-to-backscatter ratio)

$$Lr_{part}(R,\lambda) = \frac{\alpha_{part}(R,\lambda)}{\beta_{part}(R,\lambda)}$$

$$Lr_{mol}(\mathbf{k}, \boldsymbol{\lambda}) = \frac{\alpha_{mol}(R, \lambda)}{\beta_{mol}(R, \lambda)} = \frac{8\pi}{3} sr$$

dependent on range and wavelength because depends on:

- size distribution
- 🕨 shape
- composition (refractive index)

independent on range and wavelength

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In the next step:

- lidar ratio definitions are included
- ➢ the equation is reorganized
- the resulting Bernoulli equation is solved

Solution:

$$\beta_{part}(R) = -\beta_{mol}(R) + \frac{R.C.S(R) \cdot \exp\left\{-2\int_{R_0}^{R} \left[Lr_{part}(x) - Lr_{mol}\right]\beta_{mol}(x)dx\right\}}{\frac{R.C.S.(R_0)}{\beta_{part}(R_0) + \beta_{mol}(R_0)} - 2\int_{R_0}^{R} Lr_{part}(x') \cdot R.C.S.(x') \cdot \exp\left(-2\int_{R_0}^{x'} \left[Lr_{part}(x) - Lr_{mol}\right]\beta_{mol}(x)dx\right]dx'}$$

...many things ... but they are known (almost)



$\begin{aligned} & \mathcal{K}LETT-FERNALD METHOD \\ & \beta_{part}(R) = -\beta_{mol}(R) + \\ & R.C.S(R) \exp\left\{-2\int_{R_0}^{R} \left[Lr_{part}(x) - Lr_{mol}\right]\beta_{mol}(x)dx\right\} \\ & + \frac{R.C.S.(R_0)}{\beta_{part}(R_0) + \beta_{mol}(R_0)} - 2\int_{R_0}^{R} Lr_{part}(x') R.C.S.(x') \exp\left(-2\int_{R_0}^{x'} \left[Lr_{part}(x) - Lr_{mol}\right]\beta_{mol}(x)dx\right)dx' \end{aligned}$

- lidar range corrected signal (R.C.S.) is known
- molecular properties are known from meteorological data of temperature and pressure or approximated from appropriate standard atmospheres
- > particle lidar ratio is an input parameter
- \succ boundary condition at $R_0 \dots$ how to choose it?





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In principle, the solution for $\beta_{part}(R)$ can be integrated by starting from the reference range R_o , which may be either the near end (R>R_o, "forward integration") or the remote end (R<R_o, "backward integration") of the measuring range



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The most critical parameter in this method is the selection of a particle lidar ratio



Lr_{part} values independent of altitude are typically used

 $Lr_{part}(R,\lambda) \equiv \frac{\alpha_{part}(R,\lambda)}{\beta_{part}(R,\lambda)}$

Some values of particle lidar ratio:

Туре	Lr (at 532 nm)
Marine particles	20-35 sr
Saharan dust	40-70 sr
Biomass burning	70-100 sr
Urban/continental	45-75 sr

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The spectral dependence of β_{part} is strongly dependent on particle size: commonly used for the qualitative description of particle size



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In this way we obtain the profile of the particle backscatter coefficient and from this the profile of particle extinction coefficient can be estimated by:

$$\alpha_{part}(R) = Lr_{part}(R)\beta_{part}(R)$$



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Improvement : K-F method constrained by AOD_{photometer}

1 – Aplication of K-F method with an initial value of Lr_{part}

2 – Computation of aerosol extinction profile and AOD from lidar:

 $\alpha_{part}(R) = Lr_{part}\beta_{part}(R) \sum AOD_{lidar} = \int_0^\infty \alpha_{part}(R)dR = Lr_{part}\int_0^\infty \beta_{part}(R)dR$

3 – Comparison of AOD from lidar and photometer:

$$\left|AOD_{photometer} - AOD_{lidar}\right| \le \varepsilon$$

4 – Modification of Lr_{part} if necessary (go to 1)

This column-related Lr_{part} must be considered as an effective Lr_{part}: the true Lr_{part} profile remains unknown

SUMMARY ON K-F METHOD

> The K-F method allows for determining the β_{part} profile

> Numerical stability is given in the backward integration

- > The reference range R_o is usually chosen such that β_{part} at R_o is negligible compared to the known β_{mol}
- > The most critical input parameter is the $L_{part}(R)$. This quantity depends on the microphysical, chemical and morphological properties of the particles, and can vary strongly with height, specially when the atmospheric aerosols present a layered structure.
- The application of the K-F method constrained by AOD_{photometer} allows for obtaining an an effective Lr_{part}, but the true Lr_{part} profile remains unknown
- Variations between 20 and 100 sr make it practically impossible to estimate trustworthy particle extinction profiles from backscatter ones

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KOVALEV, V.A, EICHINGER, W.E. Elastic Lidar. Wiley Interscience, New Jersey, 2004. Chapters 5 (technique), 6 (uncertainties) and 7 (review on particle lidar ratio and impact).

WEITKAMP, C., Lidar. Range-resolved optical remote sensing of the Atmosphere, Springer, New York, 2005. **Chapter 7**.

As previously mentioned, the K-F method (and others that will be shown later) needs the molecular component

 β_{mol} profile can be computed from the Rayleigh theory and ideal gas law:

$$\beta_{mol}(R) = \frac{9\pi^2 (n_s^2 - 1)^2}{\lambda^4 N_s^2 (n_s^2 + 2)^2} \left(\frac{6 + 3\rho}{6 - 7\rho}\right) N_s \frac{T_0}{P_0} \frac{P(R)}{T(R)}$$

- n_s refractive index
- $\boldsymbol{\rho}$ depolarization factor (0.0301, 0.0284 and 0.0273 at 355, 532 and 1064 nm respectively)
- $N_s = 2.547 \cdot 10^{19} cm^{-3}$ molecular number density for conditions of a standard atmosphere at sea level $P_0 = 1013.25hPa$ $T_0 = 15^{\circ}C$

Only pressure and temperature profile are needed



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How to obtain T and P profiles?

Option 1: download radiosounding observations from...



Click on the image to request a sounding at that location or enter the station number above.



80222 Bogota/Eldorado (SKBO)

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How to obtain T and P profiles?

online:

Option 2: if no radiosoundings nearby, use model data from NOAA

http://www.ready.noaa.gov/READYamet.php

Air Resources Laboratory - × → C f www.ready.noaa.gov/READYamet.php Conducting research and development in the fields of air quality, atmospheric dispersion, climate, and boundary layer Enter search term(s) Go READY Archived Meteorology ARL site only O All NOAA Archived Model Graphics ARL Home Choose a forecast location by entering a 3 or 4-character station identifier or a 6-digit WMO index number or a HYSPLIT Model latitude/longitude pair and then click the Continue button, or by clicking on the location in the map. You will be taken to the model products section. Information on ARL's data archive is available at http://readv.arl.noaa.gov/archives.php. READY READY News Select a Location Transport & Dispersion Get/Run HYSPLIT Using a Code Identifier OR By Selecting a U.S. or World City ٠ Airport or WMO ID: Or choose a city --> Search for Code Volcanic Ash Short-Range Ensemble Dispersion Forecasts OR by Lat Gaussian Plume Model Latitude (degrees) nvert Deg/Min/Sec into Decimal Degrees Balloon Flight Forecasting Tools Longitude (West < 0) Current & Forecast Meteorology North America Reset Continue Archived Meteorology >> North America OR click a location on the map below. Air Quality U.S Trajectories Smoke Forecast Verification Emergency Assistance RSMC Products RSMC Information Internal Use Only READY Status READY Tools



How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA online: http://www.ready.noaa.gov/READYamet.php



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STABILITY TIME-SERIES	Choose An Archived Dataset 🔻	Go
2D MAP (NCAR GRAPHICS)	Choose An Archived Dataset •	Go
2D MAP (PSPLOT)	Choose An Archived Dataset •	Go



How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA online: http://www.ready.noaa.gov/READYamet.php



FILE FORMAT OF THE ARCHIVE DATA:

If available, the current 7 days of data are located in the file called: current7days otherwise, gdas1.mmmyy.w# where, mmm = 3 letter month (jan=January) yy = 2 number year (05=2005)



How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA

online:

http://www.ready.noaa.gov/READYamet.php

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How to obtain T and P profiles?

online:

Option 2: if no radiosoundings nearby, use model data from NOAA

http://www.ready.noaa.gov/READYamet.php

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How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA **off-line**

Download HYSPLIT software from:

https://ready.arl.noaa.gov/HYSPLIT.php

Download GDAS meteorological data from: ftp://arlftp.arlhq.noaa.gov/pub/archives/gdas1/



How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA

off-line







How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA **off-line**



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How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA

off-line X Meteorological Data Data Profile Displays a text meteorological data profile (file:profile.txt) for an ARL formatted data set. Defaults (zeros) to grid center location for the first time period. Set File Name of ARL format Data C:/hysplit4 gdas1.may14.w4 Wind Display: • Vector C Polar Time offset (hrs): • 0 0 2 0 3 0 6 0 12 0 24 0 48 Time increment (hrs): C 0 C 1 C 2 G 3 C 6 C 12 C 24 Lat: 37.16 Lon: -3.61 Profile Location Quit Help Run PROFILE



How to obtain T and P profiles?

Option 2: if no radiosoundings nearby, use model data from NOAA **off-line**

SIMULATION LOG	- • ×	
METEOROLOGICAL PROFILE LISTING Meteorological Profile: gdas1.may14.w4 File start time : 14 5 22 0 0 File ending time: 14 5 28 21 0		Radiosound
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PRSS MSLP TPP6 UMOF VMOF SHTF DSWF SOLM CSNO CICE CFZR CRAI LHTF LCLD MCLD hPa mm	RH2M U10M HCLD	comp acca
927 927 1015 0 -5E-2 -4E-2 -18.1 24.0	90.4 1.8	
Exit		
<	- 4	

Radiosoundings for the whole week each 3 h are computed

Go to folder: .../hysplit4/working to find your file "profile.txt"

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How to obtain T and P profiles?

Option 3: if no radiosoundings nearby, use standard atmosphere scaled to your surface temperature and pressure

Atmospheric	dT/dz
layer(km)	(K/km)
0-11	-6.5
11-20	0
20-32	1.0
32-47	2.8
47-51	0
51-71	-2.8
71-86	-2



Increasing

$$T = T_0 - \gamma \cdot z$$
$$p = p_0 \cdot e^{-\frac{Z}{H}}$$

where T_0 and P_0 are measured surface temperature and pressure (colocated meteorological station) and H is the so-called "scaled height" with H=8.42 km



RAMAN METHOD

To solve the problems associated with the elastic backscatter lidar several solutions have been proposed

The Raman lidar measures lidar return signals elastically backscattered by air molecules and particles and inelastically (Raman) backscattered by nitrogen and/or oxygen molecules



Raman lidar systems detect the wavelength-shifted molecular return produced by rotational and vibrational Raman scattering from the chosen molecule
RAMAN SCATTERING

Elastic scattering : scattered frequency <u>equal</u> to incident one

Raman scattering: scattered frequency <u>shifted</u> in a known amount depending on the scatterers



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RAMAN SCATTERING

Shifts for N₂ molecule:

Excitation wavelength (nm)	Shifted wavelength (nm)
355	387
488	551
532	607

The strength of Raman signals is a factor of 20 (rotational Raman lines) to 500 (vibration-rotational Raman lines) lower than the one of Rayleigh signals

The Raman lidar is mainly used during nighttime in absence of the strong daylight sky background



RAMAN LIDAR EQUATION

$$P(R,\lambda_{Raman}) = P_0(\lambda_L) \frac{C}{R^2} O(R) \beta_{Raman}(R,\lambda_{Raman}) e^{-\int_0^R [\alpha(x,\lambda_L) + \alpha(x,\lambda_{Raman})] dx}$$

 $P(R,\lambda_{Raman})$ = power received by the system after backscattering from range R

 $P_o(\lambda_L)$ = power transmitted

C = lidar calibration constant

 $\beta_{Raman}(R,\lambda_{Raman})$ (km⁻¹sr⁻¹) is the Raman molecular backscatter coefficient

 $\beta_{part}(R,\lambda_L)$ does not appear in the equation

The only particle scattering effect on the signal strength is the attenuation

 $\alpha(R,\lambda_L)$ describes the extinction on the way up to the backscatter region and $\alpha(R,\lambda_{Raman})$ the extinction on the way back to the lidar

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RAMAN METHOD: PARTICLE EXTINCTION

Starting point: Raman lidar equation

$$P(R,\lambda_{Raman}) = P_0(\lambda_L) \frac{C}{R^2} O(R) \beta_{Raman}(R,\lambda_{Raman}) e^{-\int_0^R [\alpha(x,\lambda_L) + \alpha(x,\lambda_{Raman})] dx}$$

... some considerations:

- → the overlap is assumed to be complete , O(R) =1 → $R \ge R_{min}$
- \succ β_{Raman} (R,λ_{Raman}) is calculated from the molecular number density N_{Raman} which is the N₂ or O₂ molecule number density and the Raman backscatter cross section:

$$\beta_{Raman}(R,\lambda_{Raman}) = N_{Raman}(R) \cdot \frac{d\sigma_{Raman}}{d\Omega}(\pi,\lambda_{Raman})$$

N_{Raman} (R) is calculated from actual radio-sounding observations or standard atmosphere temperature and pressure profiles

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RAMAN METHOD: PARTICLE EXTINCTION

After some operation on the Raman lidar equation that include the use of the Raman backscatter coefficient and the logarithm derivation of the equation, we obtain:

$$\alpha(R,\lambda_L) + \alpha(R,\lambda_{Raman}) = \frac{d}{dR} \ln \left[\frac{N_{Raman}(R)}{R.C.S.(R,\lambda_{Raman})} \right]$$

Total extinction coefficients are caused by particles and molecules. Therefore, the particle extinction coefficient is:

$$\begin{aligned} \alpha_{part}(R,\lambda_{L}) + \alpha_{part}(R,\lambda_{Raman}) &= \\ &= \frac{d}{dR} \ln \left[\frac{N_{Raman}(R)}{R.C.S.(R,\lambda_{Raman})} \right] - \alpha_{mol}(R,\lambda_{L}) - \alpha_{mol}(R,\lambda_{Raman}) \end{aligned}$$

RAMAN METHOD: PARTICLE EXTINCTION

To obtain the particle extinction coefficient at the transmitted wavelength we have to introduce the k-exponent which describes the wavelength dependence of the particle extinction coefficient for this spectral range:

$$\alpha_{part}(\lambda) \propto \lambda^{-k} \sum \alpha_{part}(\lambda_L) = \alpha_{part}(\lambda_{Raman}) (\lambda_L / \lambda_{Raman})^{-k}$$

 $k \approx 1$ for particles at these "small" spectral ranges [$\lambda_L, \lambda_{Raman}$] [355, 387] nm [532, 607] nm

Solution:

$$\alpha_{part}(R,\lambda_L) = \frac{\frac{d}{dR} \ln \left[\frac{N_{Raman}(R)}{R.C.S.(R,\lambda_{Raman})} \right] - \alpha_{mol}(R,\lambda_L) - \alpha_{mol}(R,\lambda_{Raman})}{1 + \left(\frac{\lambda_L}{\lambda_{Raman}} \right)^k}$$

(without any assumption on Lr_{part})

RAMAN METHOD: PARTICLE BACKSCATTER

Because now α_{part} (R) is known, $\beta_{part}(R)$ can also be calculated using the elastic lidar equation

Elastic lidar equation:

$$P(R,\lambda_L) = P_0(\lambda_L) \frac{C}{R^2} O(R) \Big[\beta_{part}(R,\lambda_L) + \beta_{mol}(R,\lambda_L) \Big] e^{-2 \int_0^R \alpha(x,\lambda_L) dx}$$

Raman lidar equation:

$$P(R,\lambda_{Raman}) = P_0(\lambda_L) \frac{C_R}{R^2} O(R) \beta_{Raman}(R,\lambda_{Raman}) e^{-\int_0^R [\alpha(x,\lambda_L) + \alpha(x,\lambda_{Raman})] dx}$$

where
$$\beta_{Raman}(R, \lambda_{Raman}) = N_{Raman}(R) \cdot \frac{d\sigma_{Raman}}{d\Omega}(\pi, \lambda_{Raman})$$

By forming the ratio: $\frac{P(R, \lambda_L)P(R_0, \lambda_{Raman})}{P(R_0, \lambda_L)P(R, \lambda_{Raman})}$

RAMAN METHOD: PARTICLE BACKSCATTER

Inserting the respective lidar equations and rearranging the resulting equation, the <u>solution</u> is:

$$\beta_{part}(R,\lambda_{L}) = -\beta_{mol}(R,\lambda_{L}) + \beta_{mol}(R_{0},\lambda_{L}) + \beta_{$$

- > overlap effects cancel out because β_{part} is determined from the ratio of two lidar equations: β_{part} is determined even at R very close to the lidar
- \succ the air density, β_{mol} and α_{mol} must be estimated from measured or standard atmosphere profiles of pressure and temperature
- Iidar signals are measured by the instrument

R

> as in the Klett method, a reference value for β_{part} at R_o must be estimated. It is recommended to choose R_o in the upper troposphere where particle scattering is negligible compared to Rayleigh scattering

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RAMAN METHOD: LIDAR RATIO

Finally, the height profile of the particle lidar ratio can be derived through the independent determination of α_{part} and β_{part} :



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SUMMARY ON RAMAN METHOD

- > The Raman method allows for independently determining the $\alpha_{part}(R)$ and $\beta_{part}(R)$, and therefore also $L_{part}(R)$
- Lidar ratio assumptions or other critical assumptions are not needed
- > The reference range R_o is usually chosen such that β_{part} at R_o is negligible compared to the known β_{mol}
- $> \beta_{part}$ can be obtained even at altitudes very close to the lidar, due to overlap effects are canceled out because β_{part} is determined from the ratio of two lidar equations
- Multiwavelength Raman lidar provides this information at several wavelengths. The "multiwavelength" information allows obtaining particle microphysical properties

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RECOMMENDED BIBLIOGRAPH

WEITKAMP, C., Lidar. Range-resolved optical remote sensing of the Atmosphere, Springer, New York, 2005. Chapter 9.

Ansmann, A., U. Wandinger, M. Riebesell, C. Weitkamp, and W. Michaelis, Independent measurement of extinction and backscatter profiles in cirrus clouds by using a combined Raman elastic-backscatter lidar, Appl. Opt., 31, 7113-7131 (1992)



- A fundamental principle of light is that the electric field E-vector of the electromagnetic wave at any instant of time displays some orientation in space
- This orientation can be fixed (linearly polarized light), rotating with time (circularly or elliptically polarized light) and randomly changing with time (natural light). Importantly, any state of polarization can be converted to any other state with the help of a set of optical devices
- > Pulsed lasers generally used in lidars produce linearly polarized light
- ➤ The lidar depolarization technique involves the transmission of a linearly polarized laser pulse and the detection via a beam splitter of the perpendicular and parallel planes of polarization of the backscattered light

- According to Mie theory, spherical particles always backscatter linearly polarized electromagnetic radiation in the same (incident) plane of polarization. A variety of approximate scattering theories predict that non-spherical particles introduce a depolarized component into the backscattered radiation
- Application of depolarization measurements is discrimination of spherical versus non-spherical aerosol particles and liquid and solid phase clouds











Granada lidar (MULHACÉN)

Starting point: elastic lidar equation

$$P(R,\lambda) = P_o(\lambda) \frac{C}{R^2} O(R) \beta(R,\lambda) e^{-2\int_0^R \alpha(x,\lambda) dx}$$

Lidars use polarized light, and many of them are able to discriminate between components. The elastic lidar equations for both components are (dropping the wavelength dependence):

$$P^{\parallel}(R) = P_o \frac{C^{\parallel}}{R^2} O(R) \beta^{\parallel}(R) e^{-2\int_0^R \alpha^{\parallel}(x) dx}$$
$$P^{\perp}(R) = P_o \frac{C^{\perp}}{R^2} O(R) \beta^{\perp}(R) e^{-\int_0^R \left[\alpha^{\parallel}(x) + \alpha^{\perp}(x)\right] dx}$$

Initial definition of "depolarization ratio":

$$"\delta(R)" = \frac{C^{\perp}\beta^{\perp}(R)e^{-\int_{0}^{R} \left[\alpha^{\parallel}(x) + \alpha^{\perp}(x)\right]dx}}{C^{\parallel}\beta^{\parallel}(R)e^{-2\int_{0}^{R} \alpha^{\parallel}(x)dx}} = \frac{\beta^{\perp}(R)}{\beta^{\parallel}(R)}e^{\int_{0}^{R} \left[\alpha^{\parallel}(x) - \alpha^{\perp}(x)\right]dx} \approx \frac{\beta^{\perp}(R)}{\beta^{\parallel}(R)}$$

Mess came from the beginning:

- different instrumental setups
- different research interests



1-. Volume linear depolarization ratio:

$$\delta_{v}(R) = \frac{\beta^{\perp}(R)}{\beta^{\parallel}(R)} = \frac{\beta_{part}^{\perp}(R) + \beta_{mol}^{\perp}(R)}{\beta_{part}^{\parallel}(R) + \beta_{mol}^{\parallel}(R)} = \frac{C^{\parallel}}{C^{\perp}} \frac{P^{\perp}(R)}{P^{\parallel}(R)} = K_{cal} \frac{P^{\perp}(R)}{P^{\parallel}(R)}$$

Features: -components: particles and molecules - signals: cross and parallel -no processing on raw data -no overlap effects



LIDAR DEPOLARIZATION Table of depolarization's parameters

Parameter	Symbol	Features	Formula
volume linear depolarization ratio	$\delta_{_{v}}$	Part.+mol ; cross and parallel; raw data	-
volume linear depolarization ratio (total)			
particle linear depolarization ratio			
particle linear depolarization ratio (total)			

 \square

LIDAR DEPOLARIZATION Table of depolarization's parameters

Parameter	Symbol	Features	Formula
volume linear depolarization ratio	δ_v	Part.+mol; cross and parallel; raw data	-
volume linear depolarization ratio (total)	$\delta_v^{\scriptscriptstyle total}$	Part.+mol ; cross and total; raw data	$\delta_{v}^{total}(R) = \frac{\delta_{v}(R)}{\delta_{v}(R) + 1}$
particle linear depolarization ratio			
particle linear depolarization ratio (total)			

3-. Particle linear depolarization ratio: $\delta_{part}(R) = \frac{\beta_{part}^{\perp}(R)}{\beta_{part}^{\parallel}(R)} = \frac{R_{backs}\delta_{v}(\delta_{mol}+1) - \delta_{mol}(\delta_{v}+1)}{R_{backs}(\delta_{mol}+1) - (\delta_{v}+1)}$

 $R_{backs} = \frac{\beta_{part} + \beta_{mol}}{\beta_{mol}} \quad \text{backscattering ratio} \\ \text{(range dependent)} \\ \delta_{mol} = \frac{\beta_{mol}^{\perp}}{\beta_{mol}^{\parallel}} \quad \text{molecular depolarization} \\ \text{(range independent)} \end{cases}$

Features:

-components: particles
-signals: cross and parallel
-retrieved backscatter coeff.
-overlap effects if K-F is used

LIDAR DEPOLARIZATION Table of depolarization's parameters

Parameter	Symbol	Features	Formula
volume linear depolarization ratio	$\delta_{_{v}}$	Part.+mol; cross and parallel; raw data	-
volume linear depolarization ratio (total)	$\delta_v^{\scriptscriptstyle total}$	Part.+mol ; cross and total; raw data	$\delta_v^{total} = \frac{\delta_v}{\delta_v + 1}$
particle linear depolarization ratio	$\delta_{\scriptscriptstyle part}$	Part. ; cross and parallel; retrieved data	$\delta_{part} = \frac{R_{backs} \delta_{v} (\delta_{mol} + 1) - \delta_{mol} (\delta_{v} + 1)}{R_{backs} (\delta_{mol} + 1) - (\delta_{v} + 1)}$
particle linear depolarization ratio (total)			

4-. Particle linear depolarization ratio (total): $\delta_{part}^{total}(R) = \frac{\beta_{part}^{\perp}(R)}{\beta_{part}^{total}(R)} = \frac{R_{backs}\delta_{v}(\delta_{mol}+1) - \delta_{mol}(\delta_{v}+1)}{(\delta_{mol}+1)(\delta_{v}+1)(R_{backs}-1)}$

Features: -components: particles -signals: cross and total -retrieved backscatter coeff. -overlap effects if K-F is used

... definitions of
$$\delta_{part}$$
 and δ_{part}^{total} :
 $\delta_{part}(R) = \frac{\beta_{part}^{\perp}(R)}{\beta_{part}^{\parallel}(R)}$
 $\delta_{part}^{total}(R) = \frac{\beta_{part}^{\perp}(R)}{\beta_{part}^{total}(R)}$
 $\delta_{part}^{total}(R) = \frac{\delta_{part}(R)}{\delta_{part}^{total}(R)}$

Summary of depolarization's parameters

Parameter	Symbol	Features	Formula
volume linear depolarization ratio	$\delta_{_{v}}$	Part.+mol; cross and parallel; raw data	_
volume linear depolarization ratio (total)	$\delta_v^{\scriptscriptstyle total}$	Part.+mol ; cross and total; raw data	$\delta_v^{total} = \frac{\delta_v}{\delta_v + 1}$
particle linear depolarization ratio	${oldsymbol{\delta}}_{\scriptscriptstyle part}$	Part. ; cross and parallel; retrieved data ,	$\delta_{part} = \frac{R_{backs}\delta_v(\delta_{mol}+1) - \delta_{mol}(\delta_v+1)}{R_{backs}(\delta_{mol}+1) - (\delta_v+1)}$
particle linear depolarization ratio (total)	$\delta^{\scriptscriptstyle total}_{\scriptscriptstyle part}$	Part. ; cross and total; retrieved data	$\delta_{part}^{total} = \frac{R_{backs}\delta_{v}(\delta_{mol}+1) - \delta_{mol}(\delta_{v}+1)}{(\delta_{mol}+1)(\delta_{v}+1)(R_{backs}-1)}$



- great range of aerosol shapes
- many aerosols consist of spherical particles: very small aerosols (small size parameter), anthropogenic aerosols, volcanic acid sulfuric droplets, sea drops released by the action of wind on water waves... All shows low volume linear depolarization ratio
- irregularly shaped aerosols (particularly volcanic and desert dusts and partially crystallized acid droplets) show larger volume linear depolarization ratio



Aerosol type	Volume linear depol. ratio
arctic haze	~ 0.013
PBL (anthropogenic)	0.02-0.10
ammonium sulfate droplets	0.02
ammonium sulfate	0.10-0.12
after crystallization	
NaCl droplets	0.06-0.12
partially crystallized	
desert dust (Sahara)	~ 0.15-0.30
desert dust (Gobbi)	~ 0.18
desert dust (Taklamakan)	0.11-0.19





Height a.s.l.(m)



SUMMARY ON LIDAR DEPOLARIZATION

- ➤ The lidar depolarization technique is based on the fact that spherical particles always backscatter linearly polarized electromagnetic radiation in the same plane of polarization, whereas non-spherical particles introduce a depolarized component into the backscattered radiation
- Due to the different definitions for depolarization parameters, it's necessary to clarify what is used (parallel vs. total; particles vs. particles+molecules)
- > Mainly, spherical particles show $\delta_v < 0.10$ and non-spherical particles show $\delta_v \sim 0.15$ -0.30

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WEITKAMP, C., Lidar. Range-resolved optical remote sensing of the Atmosphere, Springer, New York, 2005. Chapter 2 (introduction and some cloud applications).





MICROPHYSICAL PROPERTIES

In nature, particle size distributions can be described rather well by analytic expressions such as logarithmic-normal distributions:

$$dn(r) = \frac{n_{t}}{\sqrt{2\pi} \ln \sigma} \exp\left[-\frac{\left(\ln r - \ln r_{\text{mod},N}\right)^{2}}{2\left(\ln \sigma\right)^{2}}\right] d\ln r.$$

- dn(r) number concentration of particles in the radius interval $[\ln r, \ln r + d \ln r]$
 - n_t total number concentration

 $n_{\text{mod},N}$ mode radius with respect to the number concentration

 σ mode width (geometric standard deviation)

This is a monomodal distribution, however multimodal distributions (i.e. sum of \geq 2 monomodal distributions) can be found

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MICROPHYSICAL PROPERTIES

From the number concentration, other distributions can be obtained:

Number
concentration
$$dn(r) = \frac{n_t}{\sqrt{2\pi} \ln \sigma} \exp \left[-\frac{\left(\ln r - \ln r_{\text{mod},N}\right)^2}{2\left(\ln \sigma\right)^2} \right] d\ln r$$
.

Surface-area concentration $da(r) = 4\pi r^2 dn(r)$ cor distribution: dis

Integrated properties of the particle ensemble determined from the inverted distributions are:

Effective radius:

Total surface-area concentration:

Total volume concentration:

$$r_{\rm eff} = \frac{\int n(r) r^3 dr}{\int n(r) r^2 dr} \qquad a_{\rm t} = 4\pi \int n(r) r^2 dr \qquad v_{\rm t} = \frac{4\pi}{3} \int n(r) r^3 dr.$$

ume
ntration
$$dv(r) = \frac{4}{3}\pi r^3 dr^3$$

Volume
ncentration
$$dv(r) = \frac{4}{3}$$

stribution:

$$dv(r) = \frac{4}{3}\pi r^3 dn(r)$$

entration
$$dv(r) = \frac{4}{3}$$

MICROPHYSICAL PROPERTIES

A variety of methods have been proposed since the early 1970s. Examples:

Combined use of elastic lidar with other instruments

IF SUNPHOTOMETER IS USED:

- the drawback is that two colocated instruments are needed simultaneously to provide reliable data on the same particles
- different observational geometry of the instrumentation
- thus represents an additional source
- only daytime data

Inversion with regularization from Raman lidar

> - uses the spectral information contained in the backscatter and extinction profiles at several wavelengths and its change with the particle size

- limitation in number of wavel.
- in most cases only night time data
- high data quality



METHOD OF INVERSION WITH REGULATIZATION

It is the standard method for the retrieval of particle microphysical properties from multiwavelength Raman lidar observations

Input: spectrally resolved particle backscatter and extinction coefficients

The optical data are related to the physical quantities by the Fredholm integral equations of the first kind:

$$g_i(\lambda_k) = \int_{r_{\min}}^{r_{\max}} K_i(r, m, \lambda_k, s) v(r) dr + \varepsilon_i^{\exp}(\lambda_k)$$

$$i = \beta_{aer}, \alpha_{aer}, k = 1, \dots, n$$



METHOD OF INVERSION WITH REGULATIZATION $g_i(\lambda_k) = \int_{r_{min}}^{r_{max}} K_i(r, m, \lambda_k, s) v(r) dr + \varepsilon_i^{exp}(\lambda_k)$ $i = \beta_{aer}, \alpha_{aer}, \quad k = 1, ..., n$

 $g_i(\lambda_k)$ optical data at wavelengths λ_k at a specific height R

- *i* kind of information, i.e., whether it is the particle backscatter or extinction coefficient
- *k* wavelength: 355, 532 and 1064 nm

 $\mathcal{E}_i^{\exp}(\lambda_k)$ experimental error of backscatter or extinction coefficients

METHOD OF INVERSION WITH REGULATIZATION $g_i(\lambda_k) = \int_{r_{min}}^{r_{max}} K_i(r, m, \lambda_k, s) v(r) dr + \varepsilon_i^{exp}(\lambda_k)$ $i = \beta_{aer}, \alpha_{aer}, \quad k = 1, ..., n$

 $K_i(r, m, \lambda_k, s)$ kernel functions of backscatter and extinction, depending on radius, complex refractive index, wavelength and shape of particles

For spherical particles, are calculated from the respective extinction and backscatter efficiencies for individual particles weighted with their geometrical cross section πr^2

$$K_i(r, m, \lambda, s) = \frac{3}{4r} Q_i(r, m, \lambda)$$

v(r)

volume concentration distribution of particles

METHOD OF INVERSION WITH REGULATIZATION $g_i(\lambda_k) = \int_{r_{min}}^{r_{max}} K_i(r, m, \lambda_k, s) v(r) dr + \varepsilon_i^{exp}(\lambda_k)$ $i = \beta_{aer}, \alpha_{aer}, \quad k = 1, ..., n$

*r*min radius down to which particles are optically efficient. For measurements \geq 355 nm (typical in lidars), the minimum particle size is 50 nm (in radius)

 $r_{\rm max}$

radius at which concentrations are so low that particles no longer contribute significantly to the signal (typically ≤ 10 microns in troposphere)

METHOD OF INVERSION WITH REGULATIZATION

To simplify, the subscript p will be used summarizing the kind and wavelength of optical data:

$$g_{p} = \int_{r_{\min}}^{r_{\max}} K_{p}(r,m)v(r)dr + \varepsilon_{p}^{\exp} \quad \text{where} \quad p = (i, \lambda_{k})$$

- it can not be solved analytically
- the microphysical retrieval from lidars is an ill-posed inverse problem
- incompleteness of the available information: small number of wavelengths and only backscatter and extinction information is available
- non-uniqueness of the solutions: highly complex structure of tropospheric aerosols (maybe multimodal , of variable shape, particle refractive index can be wavelength-dependent)

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METHOD OF INVERSION WITH REGULATIZATION

- even uncertainties as small as round-off errors in the input data lead to disproportionately large changes in the final solution
- measurement errors are much larger than round-off errors
- different combinations of microphysical parameters may lead to similar optical properties within the measurement uncertainty
- many improvements have been done during the last decade. The most important: the reduction of measurement wavelengths to a realistic number
- the **minimum number of wavelengths is three** (355, 532 and 1064 nm) assuming simplifications for complex refractive index: the so-called 3+2 lidar systems (3 backscatters, 2 extinctions)
- the accuracy increases if backscatter up to six wavelengths are used

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To solve the equation: $g_p = \int K_p(r,m)v(r)dr + \varepsilon_p^{exp}$ where $p = (i, \lambda_k)$ $r_{\rm min}$

 \dots v(r) is discretized by a linear combination of triangular base functions $B_i(r)$ and weight factors w_i as ... $v(r) = \sum w_i B_i(r) + \varepsilon^{math}(r)$

 $B_1 B_2 B_3 B_4 B_5 B_6 B_7 B_8$

r₇

r₈

r₆

SCALE

CONCENTRATION, LINEAR

r₂

r₃

r₄ r₅

RADIUS, LOGARITHMIC SCALE

- $\varepsilon^{math}(r)$ is the mathematical residual error caused by the approximation with base functions

- $B_j(r)$ are distributed equidistantly on a logarithmic scale (to reproduce the high dinamic range of particle size distributions)
- minimum number of $B_i(r)$ is the number of input parameters
- typical number of $B_i(r)$ is 8



In general the exact position of the investigated particle size distribution over the size range used is not known

To solve this: used of a inversion window of variable width and variable position over the investigated size range

No sensible solutions are obtained if the inversion window does not cover the position of the investigated particle size distribution

Currently 50 different inversion windows within the particle size range from 0.01 to 10 μ m are used to obtain an estimate of the position of the particle size distribution

Rewriting the equations into a vector-matrix equation:



The simple solution of this equation for the weight factors is:

$$\mathbf{w} = \mathbf{A}^{-1} \mathbf{g} + \mathbf{\varepsilon}'$$



METHOD OF INVERSION WITH REGULATIZATION $w = A^{-1}g + \varepsilon'$

It fails to provide reasonable results although the optical data can be reproduced within the error limits **E**. Why?

High dynamic range of several orders of magnitude of the elements of **A** and **A**⁻¹

The term $\epsilon' = -A^{-1} \epsilon$, which describes the respective errors, and A^{-1} , which denotes the inverse of the matrix **A**, lead to error amplification of the solutions

How to solve it? Application of a procedure called regularization

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Procedure of regularization: it is used to reduce the number of solutions by restricting the highest acceptable difference between the vector Aw and **g**:

$$e^{2} \geq \left\|\varepsilon\right\|^{2} = \left\|Aw - g\right\|^{2}$$

Only solutions that minimize ε are accepted

However, due to smoothed size distributions are expected, the penalty function $\Gamma(v)$ is introduced and the minimization problem is rewritten as:

$$e^{2} \ge \left\| \varepsilon \right\|^{2} = \left\| Aw - g \right\|^{2} + \gamma \Gamma(v)$$

where $\Gamma(v) = \mathbf{w}^{T} \mathbf{H} \mathbf{w}$
itransposed
vector \mathbf{w}
$$\mathbf{w}$$



Therefore...

$$\mathbf{w} = (\mathbf{A}^{\mathrm{T}}\mathbf{A} + \gamma \mathbf{H})^{-1}\mathbf{A}^{\mathrm{T}}\mathbf{g}$$

The final solution is:

 $v(r) = \sum_{j} w_{j} B_{j}(r)$

Errors: effective radius: ~20% volume and surface-area concentrations: <±50% real part of the complex refractive index: ±0.05, imaginary part: ±50% number concentrations >±50%



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An example of forest fire smoke (17/10/2011):





METHOD OF COMBINATION OF LIDAR AND SUNPHOTOMETER

Sun-photometers: widespread tool to retrieve column integrated values of aerosol optical, microphysical and radiative properties



r the vertical profiling of aerosols in the atmosphere. roperties profiles are not easy to derive

ne of European project ACTRI Research InfraStructure Netwo

tion of AERONET sun-photon r multiwavelength lidar data profiles of aerosol microphysical properties, algorithm developed in the National Acad Belarus





LIRIC



AEROSC

NTERACTIONS

LIRIC

<u>LIDAR</u>

- 3 elastic backscattered signals (355, 532 and1064 nm)
- parallel and cross-polarized backscattered signal (532 nm)

SUN-PHOTOMETER

AERONET Retrieved properties:

- columnar particle size distribution
- volume concentrations
- refractive index
- radiative properties





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- the results of sun-photometric measurements represent the "truth" particle properties over the atmosphere
- The objective is to construct the particle vertical distribution matching both the integrated aerosol properties observed by ground-based radiometer and the vertically variable signal of multi-wavelength lidar



LIRIC





AEROSO LIDAR. $\overline{\bigcirc}$ AERONE \Box INTERACTIONS AND











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TEST #1: ZERO-BIN AND BIN-SHIFT

<u>Zero-bin test</u>: a target is placed at the output of the laser window in order to produce strong backscattered radiation. Thus, the first intense peak observed by the detector system should correspond to the zero position of our measurements



TEST #1: ZERO-BIN AND BIN-SHIFT

This test was only carried out for AN signals, since the saturation suffered in PC mode at close height range could lead to wrong results

For measuring a possible delay between AN and PC signals (bin-shift), target as clouds or aerosol layers are used



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TEST #1: ZERO-BIN AND BIN-SHIFT

Example of results for system sp-CLA-IPEN-MSP-LIDAR-I (São Paulo) $bin_0^{PC} = bin_0^{AN} + \Delta bin_{AN-PC}$

Channel (nm)	Zero-bin AN (bins)	Bin-shift (bins)	Zero-bin PC (bins)
355	8	-10	-2
387	8	-10	-2
408	8	-10	-2
532	1	-2	-1
607	6	-9	-3
660	7	-9	-2

To properly use the lidar signals, the some bins of the AN signals must be removed and several blank bins must be introduced in the PC signals following this table

TEST #2: DARK CURRENT

Dark current (DC) is the response exhibited by a receptor of radiation even during periods when it is not actively illuminated



-measurement taken by covering totally the telescope or detectors performed with enough averaging time (~10 min)

-all parameters (voltages, pulse repetition frequency ...) are configured as a usual measurement





TEST #3: TELECOVER TEST

Overlap function O(R): geometrical overlap between laser beam and telescope field of view

O(R) accounts for the partial overlap in the near height-range and tends to estabilize (ideally to 1) in the far height-range

O(R) depends, among other factors, on the alignment



Guerrero-Rascado et al., Óptica Pura y Aplicada, 2011

TEST #3: TELECOVER TEST (QUADRANTS)

<u>Telecover test (quadrants)</u>:

-based on the comparison of signals measured using different quadrants of the telescope

-named as N, E, S and W, where N is the quadrant nearest to the laser beam axis and the others named following the clockwise sense

-N measurement is performed at the end of the telecover test (N2) to check the atmospheric stability -Measurements ~2 min



TEST #3: TELECOVER TEST (QUADRANTS)

Atmospheric variability: Bad telecover!!!





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TEST #3: TELECOVER TEST (QUADRANTS)

Biaxial systems:N=N2 < E = W < SCoaxial systems:N=N2 = E = W = S



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TEST #3: TELECOVER TEST (IN-OUT)

<u>Telecover test (in-out)</u>:

-based on the comparison of signals measured using different rings of the telescope, named as IN and OUT

-IN measurement is performed at the end of the telecover test (IN2) to check the atmospheric stability

-Measurements ~2 min









TEST #4: RAYLEIGH FIT

The Rayleigh (or molecular) fit is a tool to analyze the quality of far range lidar signal:

-the range corrected signal measured by the lidar system is compared to the expected molecular range corrected signal

$$R.C.S._{mol}(z) = \beta_{mol}(z) \cdot \exp\left\{-2\int_{0}^{z} \alpha_{mol}(\xi)d\xi\right\} \equiv \beta_{mol}^{att}(z)$$

$$R.C.S.(z) = K \cdot \beta(z) \cdot \exp\left\{-2\int_0^z \alpha(\xi)d\xi\right\} \propto \beta^{att}(z)$$

-measurement ~30 min





Similar trend above 3 km agl up to 15 km agl (approx.) This height range can be used as z_{ref} for Klett-Fernald and Raman methods





AEROSO IDAR INTERACTIONS AND



AEROSOL-CLOUD INTERACTIONS

Radiative forcing by components





AERO onet and) Interactions

LIDAR APPLICATION: POLIPHON 1-STEP

 $\delta_{\rm p} = \frac{\beta_{\rm nd}^{\perp} + \beta_{\rm d}^{\perp}}{\beta_{\rm nd}^{\parallel \parallel} + \beta_{\rm d}^{\parallel}}$



As shown by Tesche et al. (2009), the particle depolarization ratio can be expressed by:

$$\delta_{\rm p} = \frac{\beta_{\rm nd} \delta_{\rm nd} (1 + \delta_{\rm d}) + \beta_{\rm d} \delta_{\rm d} (1 + \delta_{\rm nd})}{\beta_{\rm nd} (1 + \delta_{\rm d}) + \beta_{\rm d} (1 + \delta_{\rm nd})}$$

with the dust and non-dust depolarization ratios $\delta_d = 0.31$ and $\delta_{nd}=0.05$, respectively. After substituting β_{nd} with $\beta_p - \beta_d$, we solve the resulting equation to obtain a solution for β_d (for $\delta_{nd} \leq \delta_p \leq \delta_d$)

$$\beta_{\rm d} = \beta_{\rm p} \frac{(\delta_{\rm p} - \delta_{\rm nd})(1 + \delta_{\rm d})}{(\delta_{\rm d} - \delta_{\rm nd})(1 + \delta_{\rm p})}$$

 $\beta_{nd} = \beta_p - \beta_d$ For $\delta_p < \delta_{nd} \Rightarrow \beta_{nd} = \beta_p$ (contribution by dust is negligible) For $\delta_p > \delta_d \Rightarrow \beta_d = \beta_p$ (non-dust aerosol contribution is negligible)

LIDAR APPLICATION: POLIPHON 1-STEP



Illustration of the one-step and two-step POLIPHON methods to separate spherical particles from non-spherical dust particles (one-step method) and spherical particles, fine dust, and coarse dust particles (two-step method) by means of the particle depolarization ratio.



LIDAR APPLICATION: POLIPHON 2-STEP Depol. ratio [%]



(Left) 532 nm particle backscatter coefficient (green) and particle linear depolarization ratio (black), (center) particle backscatter coefficients for 1-step method, and (right) particle backscatter coefficient for 2-step method



LIDAR APPLICATION TO ACI: INC PROFILES





LIDAR APPLICATION TO ACI: INC PROFILES



Correlation between aerosol optical thickness (500 nm AOT) (i.e. (columnintegrated extinction coefficient) and column-integrated aerosol particle number concentration (column APC280) considering particles with r > 280 nm. Desert-dust-dominated observations from several field campaigns To translate the column-related findings into scales of particle extinction coefficient (measurable with lidar) and respective particle number concentration, we simply used the layer depth information from the lidar

 \Box
LIDAR APPLICATION TO ACI: ESTIMATION OF INC FROM APC280

The retrieval of APC_{280} from α_d is the basic lidar contribution to the estimation of the INC profiles

Next step is the link to the published INC parameterizations gained from comprehensive INC laboratory and field studies

The INC parameterizations introduced by DeMott et al. (2010, 2015) hold for standard (std) pressure (p_0 =1013 hPa) and temperature (T_0 =273.16 K) conditions so that we have to convert each profile value APC₂₈₀(p_z , T_z) from ambient pressure p_z and temperature T_z at height z to APC₂₈₀(p_0 , T_0) by using the factor ($T_z p_0$)/($T_0 p_z$)



LIDAR APPLICATION TO ACI: ESTIMATION OF INC FROM APC280

Global (aerosol-type-independent) INC parameterization (DeMott et al., 2010):

> $n_{\rm IN}(p_0, T_0, T_z) = a(273.16 - T_z)^b$ $\times n_{\rm a,280}(p_0, T_0)^{\left[c(273.16 - T_z) + d\right]},$

 $n_{a,280}$ (in std cm⁻³) representing APC₂₈₀ n_{IN} (in std L⁻¹) representing INC a=0.0000594 b=3.33 c=0.0265 d=0.0033temperature T(z) in K (< 273.16 K)

Finally, we transfer the obtained INC values $n_{IN}(p_0, T_0, T_z)$ to the ones for ambient pressure and temperature conditions, $n_{IN}(p_z, T_z)$, by multiplying $n_{IN}(p_0, T_0, T_z)$ with the factor $(T_0p_z)/(T_zp_0)$

Dust INC parameterization (DeMott et al., 2015):

 $n_{\rm IN}(p_0, T_0, T_z) = f_{\rm d} n_{\rm a,280}(p_0, T_0) \begin{bmatrix} a_{\rm d}(273.16 - T_z) + b_{\rm d} \end{bmatrix}$ $\times \exp[c_{\rm d}(273.16 - T_z) + d_{\rm d}],$

> $f_d=3$ $a_d=0.074$ $b_d=3.8$ $c_d=0.414$ $d_d=-9.671$



LIDAR APPLICATION TO ACI: ESTIMATION OF INC FROM APC280









European Aerosol Research Lidar Network



Aerosols, Clouds, and Trace gases Research InfraStructure Network

spalinet

Spanish and Portuguese Aerosol Lidar Network



Latin American Lidar Network



Towards operational ground based profiling with ceilometers, Doppler lidars and microwave radiometers for improving weather forecasts (COST ACTION ToProf ES1303)

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